

Kinetization of scalar field by torsion

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We show that a scalar field without a kinetic term in the Lagrangian density, coupled to a covariant divergence of the torsion vector in the Einstein-Cartan theory of gravity, becomes kinetic in its general-relativistic equivalent formulation. Dynamical scalar fields may therefore be emergent.

The Einstein-Cartan theory of gravity [1, 2] naturally extends general relativity by including the spin angular momentum of matter. The spin-spin interaction arising in this theory [3, 4] may also remove divergent integrals in quantum field theory by providing fermions with spatial extension [5] and avoid the formation of singularities from fermionic matter in black holes and in cosmology [6]. In this theory, the Lagrangian density for the gravitational field is proportional to the Ricci scalar $R = R^i_i$ [1]:

$$\mathcal{L} = -\frac{1}{2\kappa}R\sqrt{-g} + \mathcal{L}_m, \quad (1)$$

where \mathcal{L} is the total Lagrangian density, $R_{ik} = R^j_{ijk}$ is the Ricci tensor, $R^i_{mjk} = \Gamma^i_{mk,j} - \Gamma^i_{mj,k} + \Gamma^i_{lj}\Gamma^l_{mk} - \Gamma^i_{lk}\Gamma^l_{mj}$ is the curvature tensor, Γ^i_{jk} is the affine connection, the comma denotes a partial derivative with respect to the coordinates, g is the determinant of the metric tensor g_{ik} , \mathcal{L}_m is the Lagrangian density for matter, and $\kappa = 8\pi G/c^4$ is Einstein's gravitational constant. The metricity condition $g_{ij;k} = 0$, where the semicolon denotes a covariant derivative with respect to the affine connection, gives the affine connection $\Gamma^k_{ij} = \{^k_{ij}\} + C^k_{ij}$, where $\{^k_{ij}\} = (1/2)g^{km}(g_{mi,j} + g_{mj,i} - g_{ij,m})$ are the Christoffel symbols,

$$C^i_{jk} = S^i_{jk} + 2S_{(jk)}^i \quad (2)$$

is the contortion tensor, $S^i_{jk} = \Gamma^i_{[jk]}$ is the torsion tensor, $()$ denotes symmetrization, and $[]$ denotes antisymmetrization. The indices can be lowered with the metric tensor and raised with the contravariant metric tensor g^{ik} , as in general relativity [7]. The curvature tensor can be decomposed as $R^i_{klm} = P^i_{klm} + C^i_{km;l} - C^i_{kl;m} + C^j_{km}C^i_{jl} - C^j_{kl}C^i_{jm}$, where P^i_{klm} is the Riemann tensor (the curvature tensor constructed from the Christoffel symbols instead of the affine connection) and the colon denotes a covariant derivative with respect to the Christoffel symbols. We use the notation of [4]. The Ricci scalar can be decomposed as

$$R = P - 4S^i_{;i} - 4S^i S_i - C^{ijk}C_{kij}, \quad (3)$$

where $P = P^j_{ijk}g^{ik}$ is the Riemann scalar and

$$S_i = S^k_{ik} \quad (4)$$

is the torsion vector.

We consider the following Lagrangian density for matter:

$$\mathcal{L}_m = \alpha S^i \phi_{,i} \sqrt{-g}, \quad (5)$$

where ϕ is a scalar field and α is a constant. Varying the Lagrangian density with respect to the torsion tensor and equating this variation to zero gives the Cartan field equations:

$$S^j_{ik} - S_i \delta^j_k + S_k \delta^j_i = -\frac{\kappa}{2} s_{ik}^j, \quad (6)$$

where

$$s_i^{jk} = \frac{2}{\sqrt{-g}} \frac{\delta \mathcal{L}_m}{\delta C^i_{jk}} \quad (7)$$

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is the spin tensor of matter. The inverse relation is

$$S^i{}_{jk} = -\frac{\kappa}{2}(s_{jk}{}^i + \delta_{[j}^i s_{k]l}{}^l) \quad (8)$$

and its contraction gives the torsion vector:

$$S_i = \frac{1}{4}\kappa s_{ik}{}^k. \quad (9)$$

For the Lagrangian density (5), the spin tensor is

$$s_{ij}{}^k = \frac{\alpha}{2}(\delta_i^k \phi_{,j} - \delta_j^k \phi_{,i}). \quad (10)$$

Its contraction is thus

$$s_{ik}{}^k = -\frac{3\alpha}{2}\phi_{,i}. \quad (11)$$

Substituting the spin tensor and its contraction to the field equations (8) gives the torsion tensor:

$$S^i{}_{jk} = \frac{\alpha\kappa}{8}(\delta_j^i \phi_{,k} - \delta_k^i \phi_{,j}), \quad (12)$$

and the torsion vector:

$$S_i = -\frac{3\alpha\kappa}{8}\phi_{,i}. \quad (13)$$

The contortion tensor (2) is thus

$$C_{ijk} = \frac{\alpha\kappa}{4}(\phi_{,i}g_{jk} - \phi_{,j}g_{ik}). \quad (14)$$

Substituting the torsion vector and contortion tensor to (1), using (3) and (5), and omitting a covariant divergence which does not contribute to the field equations, gives

$$\mathfrak{L} = -\frac{1}{2\kappa}P\sqrt{-g} - \frac{3\alpha^2\kappa}{32}\phi_{,i}\phi_{,k}g^{ik}\sqrt{-g}. \quad (15)$$

The second term on the right-hand side of this equation is the Lagrangian density for matter in the general-relativistic equivalent formulation of the Einstein-Cartan theory. To obtain the Einstein field equations, the Lagrangian density (15) must be varied with respect to the metric tensor and such a variation must be equaled to zero.

The second term on the right-hand side of (15) has a form of a negative kinetic term for the scalar field ϕ . Such a phantom scalar field could be a source of dark energy [8]. Furthermore, the term $\alpha S^i{}_{\phi,i}\sqrt{-g}$ in the Lagrangian density (5) is dynamically equivalent (differs by a covariant divergence) to a term in which the field ϕ is coupled to the covariant divergence of the torsion vector and has a nonkinetic form:

$$\alpha S^i{}_{\phi,i}\sqrt{-g} = \alpha(S^i{}_{;i} - 2S^i S_i)\phi\sqrt{-g}. \quad (16)$$

Therefore, a kinetic scalar field can be generated from a nonkinetic scalar field by torsion. This result suggests that dynamical scalar fields may be emergent.

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