## Test of equivalence principle for particles with spin

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We consider a simple modification of the Dirac equation, such that spin-1/2 particles violate the equivalence principle, but the latter is restored by averaging over spins. An experiment is suggested to test the existence of such an effect.

It is well known<sup>1,2</sup> that, although the usual Eötvös experiment tests only the weak equivalence principle,<sup>3</sup> it could become a test for the strong equivalence principle if performed with polarized bodies (such as test bodies with aligned spins). The purpose of this paper is twofold: First we devise a theoretical model whereby spinning particles violate the equivalence principle, but the latter is restored by averaging over spins. This model involves a dimensionless coupling constant. Then we suggest an experiment which could set an upper limit on its value.

Let  $\bar{g}$  denote the local acceleration of gravity  $(g = 980 \text{ cm/sec}^2)$ . The simplest modification of the Dirac Lagrangian involving  $\bar{g}$  is to add a term proportional to  $\bar{\psi}\gamma^{\mu}g_{\mu}\psi$ . However, in a static gravitational field,  $g_0 = 0$  and  $\bar{g}$  is a gradient, therefore, such a term can be transformed away (it is similar to adding a gradient to  $A_{\mu}$ ).<sup>4</sup>

The simplest nontrivial modification of the Dirac Lagrangian is a term proportional to  $i\overline{\psi}\gamma_5\gamma^{\mu}g_{\mu}\psi$ . The Dirac equation then becomes

$$i\hbar \frac{\partial \psi}{\partial t} = (c\vec{\alpha}\cdot\vec{p} + \beta mc^2 + ik\hbar c^{-1}\gamma_5\vec{\alpha}\cdot\vec{g})\psi.$$

In the last term, which conserves CP but not Cand P separately, the factor  $\hbar c^{-1}$  has been introduced so that the "coupling constant" k is dimensionless. In the nonrelativistic limit, the additional term in the Hamiltonian is simply  $\pm k\hbar c^{-1}\overline{o} \cdot \overline{g}$ , with opposite signs for particles and antiparticles.

Such a term would mean that a spin- $\frac{1}{2}$  particle carries a gravitational dipole moment  $k\hbar c^{-1}\sigma$ . In classical language, its center of mass and center of gravity are separated by a distance  $k\hbar/mc$ . It is therefore unlikely that k is a large number.<sup>5</sup>

Let us examine the consequences of our hypothesis. First, we note that a degenerate energy level would be resolved into two close ones separated by  $2k\hbar c^{-1}g$ . As  $2\hbar c^{-1}g = 4.30 \times 10^{-23}$  eV, such a splitting would be considerably smaller than the present limit on a possible violation of the equivalence principle by weak interactions.<sup>6</sup> Moreover, the spin of a particle would precess around the vertical axis with a frequency  $2kc^{-1}g$ . Note that  $c^{-1}g = 1.03$  rad/yr. For k = 1, this precession is much too slow to be observable in neutron interference experiments.<sup>7</sup> Yet it is more than a million times faster than the one predicted by general relativity due to the dragging of inertial frames by the rotation of the earth.<sup>8</sup> However, it affects only spin, not angular momentum in general. Indeed, averaging over spins cancels the  $\overline{\sigma} \cdot \overline{g}$  term and the equivalence principle is restored on a macroscopic scale.

Consider now a *polarized* macroscopic body such as a permanent magnet. It would have an additional energy  $\pm 2kc^{-1}\overline{S} \cdot \overline{g}$  where  $\overline{S} = (\hbar/2)\sum_{i} \overline{\sigma}$  is the total spin. This induces a *torque*<sup>9</sup>  $\pm 2kc^{-1}\overline{S} \times \overline{g}$ , which could be observed in the following way:

Let the permanent magnet, thoroughly shielded from external magnetic fields, hang freely in such a way that in its equilibrium position  $\hat{S}$  is approximately horizontal. Then, if the magnetization is destroyed by heating it above the Curie point, the equilibrium position will be shifted by an angle  $\theta$ such that  $MgH\theta = 2kc^{-1}Sg$ , where M is the mass of the magnet and H the height of the point of suspension above its center of gravity. Thus,

 $\theta = 2kc^{-1}S/MH = k\hbar c^{-1}/mH,$ 

where *m* is the mass of an atom (more generally, the mass associated with spin  $\hbar/2$ ).<sup>10</sup> For iron, we get  $\theta = k(3.8 \times 10^{-16} \text{ cm/H})$ .

Unless k is very large, the main difficulties in such an experiment, apart from observing such a small angle, would be the following:

(a) External magnetic fields must be completely shielded away. Even a single quantum of magnetic flux  $\pi\hbar/e$ , spread over an area A so that  $B = \pi\hbar/eA$ , would introduce in the Hamiltonian a term similar to the one we are considering, but with a coefficient  $e\hbar B/2m_e = \pi\hbar^2/2m_eA$  instead of  $k\hbar c^{-1}g$ . We would thus need

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to be able to neglect such a term.

(b) The demagnetization process may upset the mechanical equilibrium of the test body because of the Einstein-de Haas effect  $^{11}$  and because of

<sup>1</sup>T. A. Morgan and A. Peres, Phys. Rev. Lett. <u>9</u>, 79 (1962).

<sup>2</sup>W. T. Ni, Phys. Rev. Lett. 38, 301 (1977).

<sup>3</sup>The strong equivalence principle asserts that in a freely falling, nonrotating laboratory, not only do all free particles move with constant velocities—this is the weak equivalence principle—but *all* the laws of physics are the same in that laboratory, independent of its position in space and time.

<sup>4</sup>In this paper, we made the simple assumption that the gravitational field is the gradient of a scalar field. It is well known that no scalar theory of gravitation can account for the experimental facts. However, mixed scalar-tensor theories [C. H. Brans and R. H. Dicke, Phys. Rev. <u>124</u>, 925 (1961)] or bimetric theories [N. Rosen, Ann. Phys. (N.Y.) <u>84</u>, 455 (1974)] cannot be experimentally ruled out. In such theories, in the quasistatic case, it is not difficult to construct scalars analogous to Newton's potential.

<sup>5</sup>The experimental limit on the electric dipole moment

magnetostriction.11

(1977).

It seems that the proposed experiment, although very difficult, could be feasible in the near future. *Note added in proof.* For an alternative approach to this problem, see N. D. Hari Dass, Ann. Phys. (N.Y.) 107, 337 (1977); Gen. Relativ. Gravit. 8, 89

of neutrons [W. B. Dress et al., Phys. Rev. D <u>15</u>, 9 (1977)] implies that for neutrons, k < 14 000. For

electrons, the experimental limit on k is much higher. <sup>6</sup>M. P. Haugan and C. M. Will, Phys. Rev. Lett. <u>37</u>, 1 (1976).

<sup>7</sup>R. Colella, A. W. Overhauser, and S. A. Werner, Phys. Rev. Lett. 34, 1472 (1975).

Rev. Lett. <u>34</u>, 1472 (1975). <sup>8</sup>L. I. Schiff, Phys. Rev. Lett. <u>4</u>, 215 (1960). The frequency of the Schiff precession is about  $(g/c)(R\omega/c)$ where *R* is the earth radius and  $\omega = 2\pi/day$ .

<sup>9</sup>Besides this torque, there is also a net *force* due to the gradient of  $(\vec{s} \cdot \vec{g})$ . This is, however, a much smaller effect.

<sup>10</sup>This exactly corresponds to a horizontal shift of the center of gravity by  $k\bar{h}/mc$ . We see how the principle of equivalence is violated in the present theory: A magnet suspended in an accelerated laboratory, instead of a gravitational field, would *not* tilt.

<sup>11</sup>L. D. Landau and E. M. Lifshitz, *Electrodynamics of Continuous Media* (Pergamon, Oxford, 1960).